

The Bisognano-Wichmann theorem and nets on  $\mathbf{R}^4$ .

1. Minkowski spacetime.
2. Axioms for QFT.
3. BW-theorems.
4. BW and physics.

1. Spacetime. Minkowski spacetime:  $\mathbf{R}^4$  with signature  $+- - -$ .

History: 1830s (Faraday): Concept of fields (sections of bundle over space and time). Maxwell: Field equations (set of hyperbolic PDEs on the fields). Hyperbolicity: Perturbation of one field propagates with finite speed. PDEs are invariant under the Lorentz group. Einstein: Promote the above two properties to axioms. Minkowski: 4-dimensional spacetime continuum, notion of future/past/present remains objective. Axiom: There exists a preferred class of coordinates (the inertial coordinates) such that in these coordinates the components of metric form a diagonal matrix with entries  $\pm 1$ .

The Poincaré group is defined as the group of diffeomorphisms preserving the Lorentz structure. Subgroups of the (time preserving) Lorentz group: Rotations of spacelike slices ( $SO(3)$ ) and boosts.

2. Axioms for QFT. What is physics? Naively it is a pairing  $(A, \omega) \rightarrow \omega(A)$ , where  $A$  is an observable and  $\omega$  is a state.

Pure states: We have perfect knowledge of the configuration. Assume that the set of observables is inside a  $C^*$ -algebra. Pure states are extremal points of the set of all states.

Quantum physics: Algebras become noncommutative and the standard deviation is nonzero even for pure states because of the noncommutativity.

Observable algebras and Haag-Kastler axioms.

If  $O$  is a region in the spacetime, then  $O'$  denotes its causal complement. We consider a net of  $W^*$ -algebras on the spacetime. Isotony: If  $O_1 \subset O_2$ , then  $R(O_1) \subset R(O_2)$ . Locality: If  $O_1 \subset O_2'$ , then  $R(O_1) \subset R(O_2)'$ . Poincaré symmetry:  $\alpha_g(R(O)) = R(gO)$ .

Vacuum representation. A representation of a Haag-Kastler net is called vacuum if: There is a strongly continuous representation  $U$  of  $P_+$  that implements  $\alpha_g$ :  $U(g)\pi_0(R(O))U(g)^* = \pi(\alpha_g(R(O)))$ . Also the spectrum of the generators of  $U|_{\mathbf{R}^4}$  must lie in the forward light cone. There exists a unique up to phase vacuum vector  $\Omega$  that is univariant under  $U|_{\mathbf{R}^4}$ .